

## 9 Separating Variables in the Spherical Wave Equation

There are a number of important application of PDEs in spherical coordinates. One can discuss the Heat Equation or LaPlace's Equation in (or outside) of a sphere. We will work through the Wave equation in a sphere, and leave the other (generally simpler) applications for homework exercises.

We want to consider the equation:

$$\frac{\partial^2 u}{\partial t^2} = c^2 \nabla^2 u$$

where  $u = u(\rho, \phi, \theta, t)$  and  $\nabla^2$  is the LaPlacian operator, which in Cartesian coordinates is simply:

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

(As before,  $c$  is the wave speed.)

The model for this might be something like seismic waves bouncing around inside the Earth. (However our "simple" Wave equation assumes the medium is homogeneous (thus  $c$  is a constant)—which the interior of the Earth certainly is not!)

We'll impose the Dirichlet boundary condition that on the surface of a sphere of radius  $R$ , while also implicitly assume that  $u$  is bounded in the center of the sphere and along the polar axis. In sum, the boundary conditions will be:

$$u(R, \phi, \theta, t) = 0, \quad u(\rho, \pi, \theta, t) = u(\rho, -\pi, \theta, t), \quad \frac{\partial u}{\partial \theta}(\rho, \pi, \theta, t) = \frac{\partial u}{\partial \theta}(\rho, -\pi, \theta, t)$$

as well as require  $u$  to be bounded at  $\rho = 0$ ,  $\phi = 0$ , and  $\phi = \pi$ .

To get a unique solution we'll also require the Initial conditions:

$$u(\rho, \phi, \theta, 0) = F(\rho, \phi, \theta) \quad \text{and} \quad \frac{\partial u}{\partial t}(\rho, \phi, \theta, 0) = G(\rho, \phi, \theta)$$

We begin by separating the time dependence from the spacial dependence:  $u = h(t)w(\rho, \phi, \theta)$ .

$$h''w = c^2 h \nabla^2 w \Rightarrow \frac{h''}{c^2 h} = \frac{\nabla^2 w}{w} = -\lambda$$

Thus we have:

$$\boxed{h'' + c^2 \lambda h = 0} \tag{8}$$

and

$$\nabla^2 w = -\lambda w$$

The time-dependent equation will be straight forward once we know the eigenvalues  $\lambda$ . To determine that, however, we have to treat the quite intimidating spacial equation. In spherical coordinates (where the LaPlacian much more complicated):

$$\frac{1}{\rho^2} \frac{\partial}{\partial \rho} \left[ \rho^2 \frac{\partial w}{\partial \rho} \right] + \frac{1}{\rho^2 \sin \phi} \frac{\partial}{\partial \phi} \left[ \sin \phi \frac{\partial w}{\partial \phi} \right] + \frac{1}{\rho^2 \sin^2 \phi} \frac{\partial^2 w}{\partial \theta^2} + \lambda w = 0$$

Continuing the separation by substituting:  $w = f(\rho)g(\phi)q(\theta)$ ,

$$\frac{gq}{\rho^2} \frac{d}{d\rho} \left[ \rho^2 \frac{df}{d\rho} \right] + \frac{fq}{\rho^2 \sin \phi} \frac{d}{d\phi} \left[ \sin \phi \frac{dg}{d\phi} \right] + \frac{fg}{\rho^2 \sin^2 \phi} \frac{d^2 q}{d\theta^2} + \lambda fgq = 0$$

After multiplying through by  $\left( \frac{\rho^2 \sin^2 \phi}{fgq} \right)$  we can separate off the  $\theta$  dependent terms,

$$\frac{\sin^2 \phi}{f} \frac{d}{d\rho} \left[ \rho^2 \frac{df}{d\rho} \right] + \frac{\sin \phi}{g} \frac{d}{d\phi} \left[ \sin \phi \frac{dg}{d\phi} \right] + \lambda \rho^2 \sin^2 \phi = -\frac{1}{q} \frac{d^2 q}{d\theta^2} = \mu$$

This gives a familiar eigenvalue problem for the  $\theta$  dependence:

$$\boxed{\frac{d^2 q}{d\theta^2} + \mu q = 0}, \quad q(\pi) = q(-\pi), \quad q'(\pi) = q'(-\pi) \quad (9)$$

As we saw before  $\mu = m^2$  for some non-negative integer  $m$ , while the eigenfunctions are  $\cos(m\theta)$  and  $\sin(m\theta)$ .

So that's two variables down with two to go... Unfortunately it starts to get pretty hairy from this point on. Looking at the  $\rho$  and  $\phi$  equation with  $\mu = m^2$ ,

$$\frac{\sin^2 \phi}{f} \frac{d}{d\rho} \left[ \rho^2 \frac{df}{d\rho} \right] + \frac{\sin \phi}{g} \frac{d}{d\phi} \left[ \sin \phi \frac{dg}{d\phi} \right] + \lambda \rho^2 \sin^2 \phi = m^2$$

Dividing through by  $\sin^2 \phi$ , we can separate the last two variables:

$$\frac{1}{f} \frac{d}{d\rho} \left[ \rho^2 \frac{df}{d\rho} \right] + \lambda \rho^2 = -\frac{1}{\sin \phi g} \frac{d}{d\phi} \left[ \sin \phi \frac{dg}{d\phi} \right] + \frac{m^2}{\sin^2 \phi} = \xi$$

This gives us two unfamiliar Sturm-Liouville equations:

$$\boxed{\frac{d}{d\rho} \left[ \rho^2 \frac{df}{d\rho} \right] - \xi f = -\lambda \rho^2 f} \quad (10)$$

$$(p = \rho^2, \quad q = -\xi, \quad \sigma = \rho^2)$$

$$\boxed{\frac{d}{d\phi} \left[ \sin \phi \frac{dg}{d\phi} \right] - \left( \frac{m^2}{\sin \phi} \right) g = -\xi \sin \phi g} \quad (11)$$

$$(p = \sin \phi, \quad q = -\left( \frac{m^2}{\sin \phi} \right), \quad \sigma = \sin \phi)$$

Our approach will be to analyze equation (11) first, determine the eigenvalues  $\xi$ , then substitute those eigenvalues into equation (10). The eigenvalues of equation (10) are our first separation constant  $\lambda$ , which we will use to solve the time dependent equation (8).